

THE L^2 -COHOMOLOGY OF A BOUNDED SMOOTH STEIN DOMAIN IS NOT NECESSARILY HAUSDORFF

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ABSTRACT. We give an example of a pseudoconvex domain in a complex manifold whose L^2 -Dolbeault cohomology is non-Hausdorff, yet the domain is Stein. The domain is a smoothly bounded Levi-flat domain in a two complex-dimensional compact complex manifold. The domain is biholomorphic to a product domain in \mathbb{C}^2 , hence Stein. This implies that for $q > 0$, the usual Dolbeault cohomology with respect to smooth forms vanishes in degree (p, q) . But the L^2 -Cauchy-Riemann operator on the domain does not have closed range on $(2, 1)$ -forms and consequently its L^2 -Dolbeault cohomology is not Hausdorff.

1. INTRODUCTION

For each bidegree (p, q) , with $p \geq 0, q > 0$, the Dolbeault Cohomology group $H^{p,q}(\Omega)$ of a Stein manifold Ω in degree (p, q) vanishes, and indeed this property characterizes Stein manifolds among complex manifolds (see e.g. [12, 13]). In particular, with respect to the Fréchet topology, the operator $\bar{\partial}$ from the space $\mathcal{A}^{p,q-1}(\Omega)$ of smooth $(p, q-1)$ -forms on Ω to the space $\mathcal{A}^{p,q}(\Omega)$ has closed range, since this range coincides with the null-space of the operator $\bar{\partial} : \mathcal{A}^{p,q}(\Omega) \rightarrow \mathcal{A}^{p,q+1}(\Omega)$. The aim of this paper is to show that things are much more interesting for the L^2 -cohomology of a bounded Stein domain in a Hermitian manifold.

Recall that a *Hermitian manifold* is a complex manifold whose tangent bundle has been endowed with a Hermitian metric. On a Hermitian manifold, one can define the L^2 -Dolbeault Cohomology groups, which capture the L^2 -function theory on the manifold. If the manifold Ω is realized as a relatively compact (i.e. bounded) domain inside a larger Hermitian manifold X (with the restricted metric), the L^2 -spaces of forms and functions and the L^2 -cohomology groups do not depend on the particular choice of the metric on X . According to a famous theorem of Hörmander, the L^2 -Dolbeault groups $H_{L^2}^{p,q}(\Omega)$ vanish for $q > 0$, provided Ω is a bounded Stein domain in a Stein manifold. In particular, the range of the $\bar{\partial}$ operator in the L^2 -sense is closed in each degree. One can also show that if Ω is a smoothly bounded domain with strongly pseudoconvex boundary in a complex Hermitian manifold, then again $\bar{\partial}$ has closed range, and the L^2 -Dolbeault groups are finite dimensional.

The question arises whether on a bounded Stein domain in a general complex manifold, the $\bar{\partial}$ -operator still has closed range, or equivalently whether the L^2 -cohomology groups of such a domain are Hausdorff in the natural quotient topology. This is not the case:

Theorem 1.1. *There is a compact complex surface X and a relatively compact, smoothly bounded, Stein domain \mathbb{O} in X , such that the range of the L^2 $\bar{\partial}$ -operator from the space $L_{2,0}^2(\mathbb{O})$ of square integrable $(2, 0)$ -forms to the space $L_{2,1}^2(\mathbb{O})$ of the square integrable $(2, 1)$ -forms on \mathbb{O} is not closed. Consequently, the L^2 -cohomology space $H_{L^2}^{2,1}(\mathbb{O})$ is not Hausdorff in its natural quotient topology.*

The complex surface X in the above theorem can be taken to be the product $\mathbb{P}^1 \times \mathbb{E}$, where \mathbb{P}^1 is the projective line, and \mathbb{E} is an elliptic curve defined by a rectangular lattice in \mathbb{C} . The domain \mathbb{O} defined below in Section 4 is a domain introduced by Ohsawa in [23], which has the property that in spite of the fact that

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its boundary is smooth (even real analytic), the domain \mathbb{O} is biholomorphic to a product of two planar domains, one of which is the punctured plane $\mathbb{C}^* = \mathbb{C} \setminus \{0\}$, and the other is a bounded annulus. Though \mathbb{O} is a Stein domain, its properties are quite unlike what one would expect from a bounded pseudoconvex domain in a Stein manifold. The domain \mathbb{O} used here as a counterexample, as well as closely related domains, have been studied before in several contexts (see [1, 9, 2]).

The main tool in the proof of Theorem 1.1 is the study of holomorphic extension of CR functions defined on the boundary of \mathbb{O} . On a general complex manifold M , there is no Hartogs-Bochner phenomenon (see [18] or [21]). In other words, it is not true that a CR function defined on the connected boundary $b\Omega$ of a smoothly bounded domain $\Omega \Subset M$ can be extended holomorphically into Ω . Such Hartogs-Bochner extension does take place when the ambient M is \mathbb{C}^n , $n \geq 2$, or M is Stein with dimension at least 2. In Section 3 below, we study some obstructions to holomorphic extension of CR functions, and relate them to the non-closed range property for the $\bar{\partial}$ -operator. Suppose that the range of $\bar{\partial}$ is closed in $L^2_{2,1}(\mathbb{O})$. Using the L^2 version of Serre duality (see [6]), the L^2 cohomology for $(2, 1)$ -forms $H^2_{L^2}(\mathbb{O})$ is isomorphic to $H^{0,1}_{c,L^2}(\mathbb{O})$, the L^2 cohomology for $(0, 1)$ -forms with minimal realization of the $\bar{\partial}$ -operator (see Section 2 below for definitions). One of the main observations in this paper is that if the range of $\bar{\partial}$ is closed in $L^2_{2,1}(\mathbb{O})$, then $H^{0,1}_{c,L^2}(\mathbb{O})$ vanishes. Then the $\bar{\partial}$ -Neumann operator exists on $L^2_{2,1}(\mathbb{O})$ and $H^2_{L^2}(\mathbb{O})$ vanishes. In this case, one can show that the $\bar{\partial}$ -Neumann operator is regular from W^ϵ to itself for some small $\epsilon > 0$. Then from the L^2 $\bar{\partial}$ -Cauchy problem, every CR function on the boundary $b\mathbb{O}$ in $W^{\frac{1}{2}-\epsilon}(b\mathbb{O})$ can be extended holomorphically to \mathbb{O} . On the other hand, the boundary of \mathbb{O} is Levi-flat and contains two complex tori, which divide the boundary into two disjoint parts. If we take two distinct constants on the disjoint boundary pieces, it defines a CR function on the boundary of \mathbb{O} which is in $W^{\frac{1}{2}-\epsilon}(b\mathbb{O})$, but they do not extend (see Proposition 3.1).

There are well-known elementary examples of non-Stein domains in \mathbb{C}^2 where closed range of the $\bar{\partial}$ -operator does not hold either in the Fréchet topology (cf. [25, Section 14]) or in the L^2 -topology (cf. [11, Page 75–76]). There is also an example by Malgrange of a bounded domain with pseudoconvex boundary (in a two complex dimensional torus) for which the $\bar{\partial}$ -operator does not have closed range in the Fréchet sense (cf. [22]). Similar phenomena happen in some noncompact complex Lie groups (cf. [14]). As already noted, in a Stein manifold, in the Fréchet topology on the space of smooth forms the $\bar{\partial}$ -operator has closed range. On the other hand, in the L^2 -sense, many unbounded domains in \mathbb{C}^n (for example the whole of \mathbb{C}^n itself) are easily seen not to have the closed range property for the $\bar{\partial}$ -operator. The example of \mathbb{O} considered here shows that such non-closed range phenomena are also possible on bounded Stein domains with smooth boundary. This is in sharp contrast with the case for bounded domains in \mathbb{C}^n . In fact, it has been shown recently that for a bounded Lipschitz domain in \mathbb{C}^2 such that the complement of the domain is connected in \mathbb{C}^2 , the $\bar{\partial}$ equation has closed range in $L^2_{0,1}(\Omega)$ if and only if the domain is pseudoconvex (cf. [19, Theorem 3.4]).

The range of the map $\bar{\partial} : L^2_{2,0}(\mathbb{O}) \dashrightarrow L^2_{2,1}(\mathbb{O})$ is not closed, and it is desirable to find an explicit $\bar{\partial}$ -closed L^2 $(2, 1)$ -form which is not in the L^2 range. In Section 5, we show that there is a harmonic $(2, 1)$ -form g on X such that $g|_{\mathbb{O}}$ is not in the range of $\bar{\partial}$ in the L^2 sense, but we do not know whether $g|_{\mathbb{O}}$ is in the closure of the range. On the other hand, there exists a solution v of $\bar{\partial}v = g$ which is in $W^{-\epsilon}(\mathbb{O})$ for each $\epsilon > 0$. Recalling that as a complex manifold, \mathbb{O} is the product of the punctured plane \mathbb{C}^* with an annulus, one may suspect that the lack of closed range of $\bar{\partial}$ is somehow connected to the unboundedness of the punctured plane. In Section 5, we also show that this is not the case by looking at the L^2 -function theory on \mathbb{O} in more detail, and in particular show that for $(2, 0)$ -forms which depend only on the \mathbb{C}^* factor, the $\bar{\partial}$ operator actually has closed range. This shows that the non-closed range property is more subtle.

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2. DEFINITIONS AND NOTATION

We briefly recall the definitions related to the L^2 -Dolbeault cohomology. Details on the L^2 -theory of the $\bar{\partial}$ -equation, and the associated $\bar{\partial}$ -Neumann problem may be found in standard texts on the subject, e.g. [13, 11, 7, 27], and details on the L^2 -version of Serre duality theorem used in this paper may be found in [6].

Let Ω be a Hermitian manifold, i.e., a complex manifold with a Riemannian metric which is Hermitian with respect to the complex structure on the tangent spaces. We can define in a natural way the L^2 -spaces of (p, q) -forms $L^2_{p,q}(\Omega)$. In the special case when Ω is realized as a relatively compact domain in a larger complex manifold X , and is given a Hermitian metric by restricting from X , the spaces $L^2_{p,q}(\Omega)$ are defined independently of the particular choice of the metric in X . In our application, the domain Ω will be of this latter type.

The differential operator $\bar{\partial}$ is defined classically on the space $\mathcal{A}^{*,*}(\Omega)$ of smooth forms on Ω , and maps (p, q) -forms to $(p, q + 1)$ -forms. Standard techniques of functional analysis allow us to construct extensions of the operator $\bar{\partial}$ acting as closed unbounded operators on the spaces $L^2_{p,q}(\Omega)$. Two such realizations of $\bar{\partial}$ as an unbounded operator on $L^2_{*,*}(\Omega)$ are of fundamental importance. First, the (weak) *maximal realization* of $\bar{\partial}$ is defined to have domain $\mathcal{D}_{\max}^{p,q} = \{f \in L^2_{p,q}(\Omega) \mid \bar{\partial}f \in L^2_{p,q+1}(\Omega)\}$, where $\bar{\partial}$ acts on an L^2 -form in the sense of distributions. By standard abuse of notation we will denote by $\bar{\partial}$ the unbounded, closed, densely Hilbert space operator from $L^2_{p,q}(\Omega)$ to $L^2_{p,q}(\Omega)$ with domain $\mathcal{D}_{\max}^{p,q}$, and such that for $f \in \mathcal{D}_{\max}^{p,q}$, we take $\bar{\partial}f$ in the sense of distributions. We also define the (strong) *minimal realization* $\bar{\partial}_c$ as the closed operator with smallest domain $\mathcal{D}_{\min}^{p,q}$ which extends the restriction of the $\bar{\partial}$ -operator to the smooth compactly supported forms. This may be interpreted by saying that forms in $\mathcal{D}_{\min}^{p,q}$ satisfy a boundary condition. The existence of such a “closure” $\bar{\partial}_c$ is proved in functional analysis. Note that $\mathcal{D}_{\max}^{p,q} \supsetneq \mathcal{D}_{\min}^{p,q}$ on a bounded domain Ω , and the maximal realization $\bar{\partial}$ is an extension of the operator $\bar{\partial}_c$. Note that $\bar{\partial} \circ \bar{\partial} = \bar{\partial}_c \circ \bar{\partial}_c = 0$, so we have two different examples of *Hilbert Complexes* (cf. [4]), i.e., a chain complex (in the sense of homological algebra) in which the differential operators are unbounded operators on Hilbert spaces, and the chain groups are the domains of these operators. One can associate to such a complex its cohomology, and this way we obtain, for any Hermitian manifold Ω , its L^2 -Dolbeault cohomology groups:

$$H_{L^2}^{p,q}(\Omega) = \frac{\ker(\bar{\partial} : L^2_{p,q}(\Omega) \dashrightarrow L^2_{p,q+1}(\Omega))}{\text{img}(\bar{\partial} : L^2_{p,q-1}(\Omega) \dashrightarrow L^2_{p,q}(\Omega))},$$

where the dashed arrows signify that the maximal realization is defined only on a dense subspace of the Hilbert space $L^2_{*,*}(\Omega)$ (in this case $\mathcal{D}_{\max}^{*,*}$). Similarly, we can define the L^2 -Dolbeault cohomology with *minimal realization* by setting

$$H_{c,L^2}^{p,q}(\Omega) = \frac{\ker(\bar{\partial}_c : L^2_{p,q}(\Omega) \dashrightarrow L^2_{p,q+1}(\Omega))}{\text{img}(\bar{\partial}_c : L^2_{p,q-1}(\Omega) \dashrightarrow L^2_{p,q}(\Omega))},$$

where now the minimally realized Hilbert space operators $\bar{\partial}_c$ are used. In general, the L^2 -Dolbeault groups, and the L^2 -Dolbeault groups with minimal realization are very different, but under appropriate hypotheses, there is a relation of duality between the two collections of groups, which is an analog of the classical Serre duality in the L^2 -setting (see [6] for details).

Note that the cohomology spaces $H_{L^2}^{p,q}(\Omega)$ and $H_{c,L^2}^{p,q}(\Omega)$ are complex vector spaces, and have a natural linear topology as quotients of linear topological spaces. In particular, for $H_{L^2}^{p,q}(\Omega)$, if the range $\text{img}(\bar{\partial} : L^2_{p,q-1}(\Omega) \dashrightarrow L^2_{p,q}(\Omega))$ is a closed subspace of $L^2_{p,q}(\Omega)$, then $H_{L^2}^{p,q}(\Omega)$ has the natural structure of a Hilbert space itself. However, if $\text{img}(\bar{\partial} : L^2_{p,q-1}(\Omega) \dashrightarrow L^2_{p,q}(\Omega))$ is not closed, the quotient topology is not even Hausdorff. The closed range property for the $\bar{\partial}$ -operator has other important consequences, and most importantly, it is equivalent to the possibility of solving the $\bar{\partial}$ -problem with L^2 -estimates.

The Hilbert space adjoints $\bar{\partial}^*$ and $\bar{\partial}_c^*$ are again closed, densely defined, unbounded operators on $L^2_{*,*}(\Omega)$. The *Complex Laplacian* is the operator $\square = \bar{\partial}\bar{\partial}^* + \bar{\partial}_c^*\bar{\partial}_c$, and its inverse (modulo kernel) is the $\bar{\partial}$ -Neumann

operator N . Both map the space $L^2_{p,q}(\Omega)$ to itself for each degree (p, q) . The kernel of \square in degree (p, q) consists of the *Harmonic forms*. More details on these constructions may be found in the texts mentioned above.

3. EXTENSION OF CR FUNCTIONS FROM BOUNDARIES OF MANIFOLDS

Proposition 3.1. *Let X be a Hermitian manifold of complex dimension n , and let Ω be a smoothly bounded relatively compact domain in X . Suppose that*

- (a) $b\Omega$ is connected, and there is a smooth complex hypersurface H (not necessarily connected) of X such that $H \subset b\Omega$ and $b\Omega \setminus H$ is not connected.
- (b) The L^2 -cohomology in degree $(0,1)$ with minimal realization vanishes, i.e.,

$$H_{c,L^2}^{0,1}(\Omega) = 0.$$

Then the L^2 $\bar{\partial}$ -operator from $L^2_{n,n-2}(\Omega)$ to $L^2_{n,n-1}(\Omega)$ does not have closed range. Consequently, the L^2 -cohomology group $H_{L^2}^{n,n-1}(\Omega)$ is not Hausdorff in its natural topology.

Proof. For a contradiction, assume that the $\bar{\partial}$ -operator from $L^2_{n,n-2}(\Omega)$ to $L^2_{n,n-1}(\Omega)$ has closed range. We claim that then there is an ϵ with $0 < \epsilon < \frac{1}{2}$ such that each CR function f on $b\Omega$ which belongs to the L^2 -Sobolev space $W^\epsilon(b\Omega)$ of order ϵ extends to holomorphic function F on Ω . (In the Sobolev context, “extension” means that $F \in W^{\frac{1}{2}+\epsilon}(\Omega)$, and the Sobolev trace of F on $b\Omega$ is f). Postponing the proof of the claim for now, to produce a contradiction it suffices to produce a CR function on $b\Omega$ which is of class W^s on $b\Omega$, for each $0 < s < \frac{1}{2}$, and which does not admit a holomorphic extension to Ω . But such a function is easy to construct as follows.

Let $b\Omega^+$ be a connected component of $b\Omega \setminus H$, and let $b\Omega^-$ be the union of the remaining connected components. By hypothesis, neither of $b\Omega^\pm$ is empty, their union is $b\Omega \setminus H$, and the closures of $b\Omega^+$ and $b\Omega^-$ meet along the complex hypersurface H . Define a locally integrable function f on $b\Omega$ by setting $f \equiv 1$ on $b\Omega^+$ and $f \equiv 0$ on $b\Omega^-$.

We claim that f is CR on $b\Omega$ in the sense of distributions. This is clear except along the complex hypersurface H , where f has a jump discontinuity. To verify the fact that f is CR on H , let $p \in H$, and choose a local C^∞ real coordinate system $(x_1, y_1, \dots, x_{n-1}, y_{n-1}, t)$ on a neighborhood of p in $b\Omega$ such that the hypersurface H is given by $\{t = 0\}$, and along H , $z_j = x_j + iy_j$ is a complex coordinate on H for $1 \leq j \leq n-1$. We can further assume that $f = 1$ on $\{t > 0\}$ and $f = 0$ on $\{t < 0\}$. At the point p , a basis of $(0,1)$ -vector fields is given by the $(n-1)$ vectors $\frac{\partial}{\partial \bar{z}_j}$, where $1 \leq j \leq n-1$. It follows that in a neighborhood of p in $b\Omega$ there are vector fields

$$Z_j = \frac{\partial}{\partial \bar{z}_j} + \sum_{\ell=1}^{n-1} \left(a_j^\ell \frac{\partial}{\partial x_\ell} + b_j^\ell \frac{\partial}{\partial y_\ell} \right) + c \frac{\partial}{\partial t},$$

which pointwise span the CR vector fields, where a_j^ℓ, b_j^ℓ and c (with $1 \leq j, \ell \leq n-1$) are smooth functions near p which vanish identically on the hypersurface H . To show that f is CR, we need to check that for each j , we have $Z_j f = 0$ in the sense of distributions (in order to induce a distribution corresponding to the function f , we use the standard volume form of \mathbb{R}^{2n-1}). For each ϕ which is smooth and compactly supported near p , we need to show that $(Z_j f)(\phi) = 0$, i.e., $\int_{t>0} Z_j^t \phi = 0$, where $Z_j^t = -\frac{\partial}{\partial \bar{z}_j} + Y_j + \lambda_j$,

where Y_j is a smooth vector field and λ_j a smooth function given by

$$Y_j = - \left(\sum_{\ell=1}^{n-1} \left(a_j^\ell \frac{\partial}{\partial x_\ell} + b_j^\ell \frac{\partial}{\partial y_\ell} \right) + c \frac{\partial}{\partial t} \right) \quad (3.1)$$

$$\lambda_j = - \left(\sum_{\ell=1}^{n-1} \left(\frac{\partial a_j^\ell}{\partial x_\ell} + \frac{\partial b_j^\ell}{\partial y_\ell} \right) + \frac{\partial c}{\partial t} \right). \quad (3.2)$$

Now, assuming that the support of ϕ is contained in the cube $\{|x_j| < \delta, |y_j| < \delta, |t| < \delta\}$, we have using Fubini's theorem (the hat $\widehat{}$ in the formulas below indicates that this factor is absent from the product):

$$\begin{aligned} \int_{t>0} \frac{\partial \phi}{\partial \bar{z}_j} &= \int \left(\int_{|x_j|<\delta, |y_j|<\delta} \frac{\partial \phi}{\partial \bar{z}_j} dx_j dy_j \right) dx_1 dy_1 \dots \widehat{dx_j} \widehat{dy_j} \dots dx_{n-1} dy_{n-1} dt \\ &= \int_{t=0}^{\delta} 0 \cdot dx_1 dy_1 \dots \widehat{dx_j} \widehat{dy_j} \dots dx_{n-1} dy_{n-1} dt \\ &= 0, \end{aligned}$$

where the inner integral in the repeated integral vanishes by an application of the divergence formula. To complete the proof we need to show that $\int_{t>0} (Y_j + \lambda_j)\phi = 0$. Now let ψ be a compactly supported smooth function on the real line which is identically 1 in a neighborhood of 0, and for small $\epsilon > 0$, let $\phi_\epsilon(x_1, y_1, \dots, x_{n-1}, y_{n-1}, t) = \phi(x_1, y_1, \dots, x_{n-1}, y_{n-1}, t)\psi(\frac{t}{\epsilon})$. Writing $\phi = \phi_\epsilon + (\phi - \phi_\epsilon)$, we see that $\int_{t>0} (Y_j + \lambda_j)(\phi - \phi_\epsilon) = 0$, since f is a smooth CR function outside H . The integral $\int_{t>0} \lambda_j \phi_\epsilon$ is clearly $O(\epsilon)$ as $\epsilon \rightarrow 0$, since the function $\lambda_j \phi_\epsilon$ is bounded uniformly in ϵ , and the integral ranges over a subset of the support of ϕ_ϵ , which has volume $O(\epsilon)$. Noting that the coefficients a_j^ℓ, b_j^ℓ, c vanish along H , we see that they are $O(\epsilon)$ in the support of ϕ_ϵ . Further, $\left| \frac{\partial}{\partial t}(\phi_\epsilon) \right| = O\left(\frac{1}{\epsilon}\right)$. Therefore,

$$\begin{aligned} \left| \int_{t>0} Y_j \phi_\epsilon \right| &\leq \sup \left| \sum_{\ell=1}^{n-1} \left(a_j^\ell \frac{\partial \phi_\epsilon}{\partial x_\ell} + b_j^\ell \frac{\partial \phi_\epsilon}{\partial y_\ell} \right) + c \frac{\partial \phi_\epsilon}{\partial t} \right| \cdot \text{vol}(\text{support } \phi_\epsilon) \\ &\leq C(\epsilon + \epsilon + \epsilon \cdot \frac{1}{\epsilon}) \cdot \epsilon \\ &= O(\epsilon). \end{aligned}$$

Combining the estimates, it follows that for each $\epsilon > 0$, we have that

$$\begin{aligned} |(Z_j f)(\phi)| &= \left| \int_{t>0} Z_j^t \phi \right| \\ &= \left| \int_{t>0} -\frac{\partial \phi}{\partial \bar{z}_j} + Y_j \phi + \lambda_j \phi \right| \\ &\leq C\epsilon, \end{aligned}$$

so that $Z_j f = 0$ for $1 \leq j \leq n-1$, i.e., f is CR in the sense of distributions.

We now claim that for $0 < s < \frac{1}{2}$, the CR function f belongs to the fractional-order Sobolev space $W^s(b\Omega)$. Let B denote the unit ball of \mathbb{R}^{2n-1} and let B^\pm denote the upper and the lower half balls. To see that $f \in W^s(b\Omega)$, it clearly suffices to show $u \in W^s(B)$, where u is the function on B which is identically 1 on B^+ and which vanishes on B^- . The fact that $u \in W^s(B)$ for $0 < s < \frac{1}{2}$ can be verified by a direct computation, cf. [28, Proposition 5.3], or it follows from [20, Theorem 11.4], where it is shown that the extension by 0 of a function in $W^s(B^+)$ is a function in $W^s(B)$ if $0 < s < \frac{1}{2}$.

Note that the function f does not admit a holomorphic extension to either Ω or $X \setminus \Omega$. If such an extension F did exist, by standard estimates on Bochner-Martinelli type singular integrals, F would be C^∞ -smooth up to the boundary on $b\Omega \setminus H$. Then by a classical boundary uniqueness result, the holomorphic

function F has to be simultaneously both identically 1 and identically 0 on Ω . This produces the required contradiction. To complete the proof we only need to establish the claim made in the first paragraph regarding the existence of an ϵ with $0 < \epsilon < \frac{1}{2}$ for which each CR function of class $W^\epsilon(b\Omega)$ extends holomorphically into Ω .

By assumption, the $\bar{\partial}$ -operator has closed range in $L^2_{n,n-1}(\Omega)$. Then the $\bar{\partial}$ -Neumann operator $\mathbf{N} = \mathbf{N}_{n,n-1}$ exists on Ω as a bounded operator on $L^2_{n,n-1}(\Omega)$, since in the top degree, $\bar{\partial}$ automatically has closed range from $L^2_{n,n-1}(\Omega)$ to $L^2_{n,n}(\Omega)$. Further, by L^2 -Serre duality (cf. [6, Theorem 2]) it follows that $H^2_{L^2, n-1}(\Omega) = 0$, and further that the harmonic space $\mathcal{H}^{n,n-1}(\Omega) = 0$ (cf. [6]). Using an observation of Kohn (see Proposition 3.2 below), there is an $0 < s < \frac{1}{2}$ such that the operator \mathbf{N} extends as a bounded operator on $W^s_{n,n-1}(\Omega)$, i.e., the space of $(n, n-1)$ -forms on Ω with coefficients in $W^s(\Omega)$. When \mathbf{N} maps W^s to itself, it follows that the ‘‘canonical solution operator’’ $\bar{\partial}^* \mathbf{N}$ and all the related operators ($\bar{\partial} \mathbf{N}$ and the Bergman projection) also map W^s to itself. (For integral s , a proof based on a standard commutator estimate may be found in texts on the $\bar{\partial}$ -Neumann problem, see [7, Theorem 6.1.4 and Theorem 6.2.2] or [27, Lemma 3.2 and Corollary 3.3]. Essentially the same arguments generalize to fractional s , even with $0 < s < 1$, if we replace the tangential differential operators in the commutator estimates by tangential pseudo-differential operators of fractional order and use the commutator estimates for such operators, see [11, Appendix, Prop. A.5.1 to A.5.3]). Therefore $\bar{\partial} \mathbf{N}$ and $\bar{\partial}^* \mathbf{N}$ map $W^s(\Omega)$ to itself. If $0 < s < \frac{1}{2}$, the space $\mathcal{C}_0^\infty(\Omega)$ is dense in $W^s(\Omega)$, so that $W^s(\Omega)$ coincides with $W^s_0(\Omega)$ (cf. [20, Theorem 11.1]), and therefore $W^{-s}(\Omega)$, defined to be the dual of $W^s_0(\Omega)$, is actually the dual of $W^s(\Omega)$. Since \mathbf{N} is self-adjoint, it maps $W^{-s}(\Omega)$ to itself. It follows that the associated operator $\bar{\partial} \mathbf{N}$ (or $\bar{\partial}^* \mathbf{N}$) maps $W^{-s}(\Omega)$ to itself.

Now let f be a CR function on $b\Omega$ belonging to $W^\epsilon(b\Omega)$ where $\epsilon = \frac{1}{2} - s$. We can extend f to a function $\tilde{f} \in W^{\frac{1}{2}+\epsilon}(\Omega)$. Then $\bar{\partial} \tilde{f} \in W^{\epsilon-\frac{1}{2}}(\Omega) = W^{-s}(\Omega)$. We consider the function

$$u_c = - * \bar{\partial} \mathbf{N} (* \bar{\partial} \tilde{f}),$$

where $*$ denotes the Hodge Star operator on the space of differential forms on the Hermitian manifold Ω . Since $*$ induces an isometry of each Sobolev space, it follows that $u_c \in W^{-s}(\Omega)$. It follows from [6, Theorem 3], that u_c satisfies $\bar{\partial}_c u_c = \bar{\partial} \tilde{f}$, where $\bar{\partial}_c$ is the minimal realization of the $\bar{\partial}$ operator (see [6] or the proof of Theorem 3.2 in [26]). Furthermore, using the regularity of $\bar{\partial} \mathbf{N}$ on $W^{-s}(\Omega)$ established above, we have

$$\|u_c\|_{W^{-s}(\Omega)} \leq C \|\bar{\partial} \tilde{f}\|_{W^{-s}(\Omega)} \leq C \|\tilde{f}\|_{W^{1-s}(\Omega)} \leq C \|f\|_{W^{\frac{1}{2}-s}(b\Omega)}.$$

If we define

$$F = \tilde{f} - u_c,$$

then $F \in W^{-s}(\Omega)$ and F is holomorphic. Using the weak and strong extension of $\bar{\partial}_c$ (cf. [19, Lemma 2.4]), we see that u_c satisfies the equation $\bar{\partial} u_c = \bar{\partial} \tilde{f}$ in an open neighbourhood of $\bar{\Omega}$ if we set u_c and $\bar{\partial} \tilde{f}$ equal to zero outside Ω . From the interior regularity for $\bar{\partial}$, we have that u_c is in $W^{1-s}(\bar{\Omega})$. This shows that u_c actually has trace in $W^{\frac{1}{2}-s}(b\Omega)$. The trace of u_c on $b\Omega$ is equal to zero since u_c satisfies the $\bar{\partial}$ -Cauchy problem. Noting that the distributional boundary value of F on $b\Omega$ is of class $W^{\frac{1}{2}-s}(b\Omega)$, it follows that $F \in W^{1-s}(\Omega)$, so that F is indeed a holomorphic extension of f . \square

To complete the argument in the previous result, we will use the following observation, which is due to Kohn (see [3]).

Proposition 3.2. *Let Ω be a relatively compact domain with smooth boundary in a Hermitian manifold X . Suppose that for some degree (p, q) , the $\bar{\partial}$ -Neumann operator $\mathbf{N}_{p,q}$ exists on Ω as a bounded linear operator on the space $L^2_{p,q}(\Omega)$ of square integrable (p, q) -forms, and suppose that the harmonic space of degree (p, q) is trivial. Then there is an ϵ_0 , with $0 < \epsilon_0 \leq \frac{1}{2}$, such that for $0 < \epsilon < \epsilon_0$, the operator $\mathbf{N}_{p,q}$ extends as a bounded operator from the Sobolev space $W^\epsilon_{(p,q)}(\Omega)$ to itself.*

If Ω admits a bounded Hölder continuous plurisubharmonic exhaustion (which happens when $X = \mathbb{C}^n$ [8] or X is a complex projective space [24, 5]) then the conclusion follows directly. However, the existence of such a bounded plurisubharmonic exhaustion is not necessary for the conclusion to hold. In our application, we will use it on a domain which does not admit a bounded plurisubharmonic exhaustion (see Lemma 4.3). The proof of the proposition is the same as the arguments used in Kohn-Nirenberg [16] using pseudo-differential operators (see [15]) and we omit the details.

4. THE PROOF OF THEOREM 1.1

4.1. The domain \mathbb{O} . As was already mentioned in the introduction, the domain \mathbb{O} will be realized as a domain with real-analytic boundary in the compact complex surface $X = \mathbb{P}^1 \times \mathbb{E}$, where \mathbb{P}^1 is the complex projective line (the Riemann Sphere) and \mathbb{E} is an elliptic curve defined in the following way. Let $\beta > 0$, and let $\Gamma_\beta = \{e^{2k\pi\beta} | k \in \mathbb{Z}\}$ be the cyclic subgroup of $\mathbb{C}^* = \mathbb{C} \setminus \{0\}$ generated by the number $e^{2\pi\beta}$. Then \mathbb{E} is the quotient $\mathbb{C}^*/\Gamma_\beta$ which is an elliptic curve.

Let the natural projection $\mathbb{C}^* \rightarrow \mathbb{E}$ be denoted by $w \mapsto [w]$ (so that $[e^{2\pi\beta}w] = [w]$) and let $z : \mathbb{P} \rightarrow \mathbb{C} \cup \{\infty\}$ denote the inhomogeneous coordinate on \mathbb{P} . Then \mathbb{O} is the domain in $\mathbb{P} \times \mathbb{E}$ given by

$$\mathbb{O} = \{(z, [w]) \in \mathbb{P} \times \mathbb{E} : \operatorname{Re}(zw) > 0\}, \quad (4.1)$$

where it is easily seen that the condition $\operatorname{Re}(zw) > 0$ is well defined independently of the choice of the lift w of the point $[w] \in \mathbb{E}$. This is a smoothly bounded Levi-flat domain in $X = \mathbb{P} \times \mathbb{E}$, and is in fact biholomorphic to a product domain in \mathbb{C}^2 , where one factor is \mathbb{C}^* and the other is an annulus. Indeed, let

$$\mathbb{A} = \left\{ W \in \mathbb{C} : e^{-\frac{\pi}{2\beta}} < |W| < e^{\frac{\pi}{2\beta}} \right\}, \quad (4.2)$$

which is an annulus in the plane. Let Φ be the map from $\mathbb{C}^* \times \mathbb{A}$ to X given by

$$\Phi(Z, W) = (Z, [Z^{-1} \exp(i\beta \cdot \log W)]) = (Z, [Z^{-1} \cdot W^{i\beta}]). \quad (4.3)$$

Φ is well-defined in spite of the multivaluedness of the logarithm. It is not difficult to verify that Φ is a biholomorphism from $\mathbb{C}^* \times \mathbb{A}$ onto \mathbb{O} . We will refer to $\Omega = \mathbb{C}^* \times \mathbb{A}$ as the *product model* of \mathbb{O} .

Note that the domain \mathbb{O} depends on the choice of the parameter β , and the domains obtained for distinct β are easily seen to be non-biholomorphic. Therefore, in fact we have a one-parameter family of counterexamples to prove Theorem 1.1. In the sequel, we consider the Ohsawa domain corresponding to one fixed β .

The Levi structure of $b\mathbb{O}$ can be summarized as follows:

Proposition 4.1. *$b\mathbb{O}$ is a smooth, real-analytic, connected, Levi-flat hypersurface. The complex tori $\{0\} \times \mathbb{E}$ and $\{\infty\} \times \mathbb{E}$ are contained in $b\mathbb{O}$ and the complement of these two tori is a disjoint union of two open subsets Σ^\pm of $b\mathbb{O}$. Each of Σ^\pm is CR equivalent to the product $\mathbb{C}^* \times S^1$ (with the natural CR structure).*

Proof. The assertions in the first sentence can be verified by direct computation, starting from the representation (4.1).

The map Φ extends biholomorphically to a neighborhood of $\mathbb{C}^* \times \overline{\mathbb{A}}$, and it is easy to see that the image of $\mathbb{C}^* \times b\mathbb{A}$ is all of $b\mathbb{O}$ except the two tori $\{0\} \times \mathbb{E}$ and $\{\infty\} \times \mathbb{E}$. Note that the boundary of the annulus \mathbb{A} consists of two circles $b\mathbb{A}^+$ and $b\mathbb{A}^-$. Define $\Sigma^+ = \Phi(\mathbb{C}^* \times b\mathbb{A}^+)$ and $\Sigma^- = \Phi(\mathbb{C}^* \times b\mathbb{A}^-)$. Then $b\mathbb{O}$ is the disjoint union of the four pieces Σ^+ , Σ^- , $\{0\} \times \mathbb{E}$ and $\{\infty\} \times \mathbb{E}$. By construction each of Σ^\pm is a Levi-flat hypersurface biholomorphically equivalent to $\mathbb{C}^* \times S^1$ (where S^1 is the circle). \square

If we use (Z, W) as coordinates on \mathbb{O} , each function holomorphic on \mathbb{O} has a Laurent expansion

$$\sum_{(j,k) \in \mathbb{Z}^2} a_{jk} Z^j W^k. \quad (4.4)$$

Viewed as a complex manifold, being a product of planar domains, \mathbb{O} is Stein, and the Dolbeault Cohomology $H^{p,q}(\mathbb{O})$ vanishes for each positive q . Note further there are nonconstant bounded holomorphic

functions on \mathbb{O} , namely, the ones represented in the product model as bounded holomorphic functions of W alone.

In order to study L^2 theory on \mathbb{O} , we need to impose an arbitrary Hermitian metric on X and restrict it to \mathbb{O} . The actual L^2 -spaces of forms and functions, the realizations of the $\bar{\partial}$ -operator are independent of the choice of the metric. For simplicity therefore we give the most symmetric metric to X , which arises as the product metric of constant curvature metrics on the factors. We endow \mathbb{P} (which is diffeomorphic to the round 2-sphere) with a round metric (the Fubini-Study metric) and the elliptic curve \mathbb{E} with a flat metric. The metric on \mathbb{P} is normalized such that

$$\phi = \frac{dz}{1 + |z|^2}, \quad (4.5)$$

is a $(1,0)$ -form of unit length at each point. Denote by ψ the unique $(1,0)$ -form on \mathbb{E} whose pullback to \mathbb{C}^* by the map $\pi : w \mapsto [w]$ satisfies

$$\pi^* \psi = \frac{dw}{w}. \quad (4.6)$$

Such a ψ exists since the form $\frac{dw}{w}$ is periodic with respect to the action of the group Γ_β by multiplication on \mathbb{C}^* . Then ψ is in fact a holomorphic 1-form on \mathbb{E} , and by declaring it to be of unit length we get a flat Hermitian metric on \mathbb{E} . Therefore, the metric on \mathbb{O} induced from $\mathbb{P} \times \mathbb{E}$ is represented as

$$ds^2 = \phi \otimes \bar{\phi} + \psi \otimes \bar{\psi}, \quad (4.7)$$

where, by standard abuse of notation, we denote also by ϕ and ψ the pullbacks of these forms from \mathbb{P} and \mathbb{E} respectively to the product $\mathbb{P} \times \mathbb{E}$ via projections on the factors.

We can pull back the metric (4.7) on \mathbb{O} via the map Φ and obtain a metric $\Phi^*(ds^2)$ on the product model $\Omega = \mathbb{C}^* \times \mathbb{A}$, so that with these metrics, Φ becomes an isometry. By a direct computation we can verify that the Riemannian volume form on Ω associated to the pullback metric $\Phi^*(ds^2)$ is represented as:

$$\omega = \frac{\beta^2}{(1 + |Z|^2)^2 |W|^2} dV \quad (4.8)$$

where $dV = (-2i)^{-2} dZ \wedge d\bar{Z} \wedge dW \wedge d\bar{W}$ represents the Euclidean volume form on \mathbb{C}^2 . We characterize some Sobolev spaces of holomorphic functions:

Lemma 4.2. (1) *The Bergman space $\mathcal{O}(\mathbb{O}) \cap L^2(\mathbb{O})$ consists of functions which are represented in the product model $\mathbb{C}^* \times \mathbb{A}$ as functions on \mathbb{A} alone, with no dependence on \mathbb{C}^* .*

(2) *The Holomorphic Sobolev space $\mathcal{O}(\mathbb{O}) \cap W^1(\mathbb{O})$ consists of only the constant functions.*

(3) *The above statements remain true if \mathbb{O} is replaced by $X \setminus \bar{\mathbb{O}}$.*

Proof. The proofs are by direct computations. For (1), it suffices to note that in the product representation, thanks to the structure of the volume form (4.8), the terms of the Laurent series (4.4) are orthogonal. Consequently, it suffices to prove that terms of the form $Z^j W^k$ are not in L^2 if $j \neq 0$. But this follows from integration on the product representation $\mathbb{C}^* \times \mathbb{A}$ using the volume form (4.8).

Thanks to part (1) it suffices to show that the functions W^k , where $k \neq 0$ are not in the Sobolev space $W^1(\mathbb{O})$. To show this, it suffices to show that the differential $d(W^k) = kW^{k-1}dW$ is not in the space $L^2_{1,0}(\mathbb{O})$. Since $W^{i\beta} = zw$, differentiating logarithmically, we see that

$$\begin{aligned} i\beta \frac{dW}{W} &= \frac{dw}{w} + \frac{dz}{z} \\ &= \psi + \frac{1 + |z|^2}{z} \phi, \end{aligned}$$

where ψ and ϕ are as in (4.6) and (4.5) respectively. Using $z = Z$ and the pointwise orthogonality of ψ and ϕ , we have for the pointwise norm:

$$\beta^2 \frac{|dW|^2}{|W|^2} = 1 + \frac{(1 + |Z|^2)^2}{|Z|^2}.$$

Using the volume form (4.8), we now see that $W^{k-1}dW$ is not square integrable for any k , so that it follows that the function W^k is not in the Sobolev space $W^1(\mathbb{O})$. Therefore, the only functions in $\mathcal{O}(\mathbb{O}) \cap W^1(\mathbb{O})$ are the constants.

The last statement follows since \mathbb{O} and $X \setminus \overline{\mathbb{O}}$ are isometrically biholomorphic by the map $(z, w) \rightarrow (-z, w)$. □

The following lemma shows that in spite of the boundedness of the domain \mathbb{O} in $\mathbb{P} \times \mathbb{E}$, in some respects the function theory on \mathbb{O} is analogous to that on an unbounded domain:

Lemma 4.3. (a) *Bounded plurisubharmonic functions on \mathbb{O} do not separate points. In particular, \mathbb{O} does not admit a bounded plurisubharmonic exhaustion function.*

(b) *A continuous plurisubharmonic function on \mathbb{O} is a constant.*

Proof. For (a), representing \mathbb{O} as a product, we see that any bounded plurisubharmonic function on \mathbb{O} must be constant on each slice $\mathbb{C}^* \times \{w\}$ for each $w \in \mathbb{A}$.

For (b), note that the restriction of such a continuous plurisubharmonic function to the torus $T_0 = \{0\} \times \mathbb{E}$ is a constant. But it is easy to see that T_0 is contained in the closure of each slice $\mathbb{C}^* \times \{w\}$ of the product representation. Consequently, the constant value assumed by the plurisubharmonic function on each slice is the same. □

Remark: The function $\operatorname{Re}(zw)$ is pluriharmonic (hence plurisubharmonic), and serves as a defining function for the domain \mathbb{O} except near the torus $T_\infty = \{\infty\} \times \mathbb{E}$. But it is not a global defining function of \mathbb{O} since it is not defined near T_∞ . A related result on the non-existence of a plurisubharmonic exhaustion function on a certain Stein domain with Levi-flat boundary may be found in [24, Theorem 1.2].

4.2. Proof of Theorem 1.1. We can now complete the proof of Theorem 1.1. Let

$$H = (\{0\} \times \mathbb{E}) \cup (\{\infty\} \times \mathbb{E}).$$

It is easy to see that $b\mathbb{O}$ is actually connected, and $b\mathbb{O} \setminus H$ is not connected, since it is the disjoint union of Σ^+ and Σ^- .

To complete the proof of Theorem 1.1 we will study the cohomology group $H_{c,L^2}^{0,1}(\mathbb{O})$. This L^2 -cohomology group is independent of the choice of the metric adopted, so we can simplify our work by choosing the metric (4.7). It suffices to endow the ambient manifold $X = \mathbb{P} \times \mathbb{E}$ with a Hermitian metric, which then can be restricted to the domain \mathbb{O} . We now prove the following:

Lemma 4.4. *With the metric (4.7) (and therefore with any other comparable metric), suppose that the range of $\bar{\partial} : L_{2,0}^2(\mathbb{O}) \rightarrow L_{2,1}^2(\mathbb{O})$ is closed in $L_{2,1}^2(\mathbb{O})$. Then we have*

$$H_{c,L^2}^{0,1}(\mathbb{O}) = 0$$

Proof. Let f be a $\bar{\partial}_c$ -closed $(0,1)$ -form on \mathbb{O} . We need to show that the equation $\bar{\partial}_c u = f$ has a solution. We instead first consider the equation on $X = \mathbb{P} \times \mathbb{E}$ given as $\bar{\partial} \tilde{u} = \tilde{f}$, where \tilde{f} is the $\bar{\partial}$ -closed $(0,1)$ -form obtained by extending the form f as 0 outside \mathbb{O} . By the Hodge decomposition on the compact Kähler manifold X , this equation has a solution provided \tilde{f} is orthogonal to the Harmonic space $\mathcal{H}^{0,1}(X)$. But by the Künneth formula, $\mathcal{H}^{0,1}(X)$ is one dimensional and generated by the form $\bar{\psi}$, where ψ is as in (4.6), and so there is an \tilde{u} satisfying the equation provided $(\tilde{f}, \bar{\psi})_X = 0$.

Define a smooth (0,2)-form on \mathbb{O} by setting $v = z\bar{\phi} \wedge \bar{\psi}$. One easily sees that v is not in $L^2_{0,2}(\mathbb{O})$. Recall that, if ϑ denotes the formal adjoint of the differential operator $\bar{\partial}$, we can write $\vartheta = -*\bar{\partial}*$, where $*$ denotes the Hodge star operator on the de Rham complex of X , which is a \mathbb{C} -antilinear map (see [6] for details), and we have

$$\begin{aligned} \vartheta v &= -*\bar{\partial}*(z\bar{\phi} \wedge \bar{\psi}) \\ &= *\bar{\partial}(\bar{z}\phi \wedge \psi) \\ &= *\bar{\partial}\left(\frac{\bar{z}}{1+|z|^2}dz \wedge \psi\right) \\ &= *\left(\frac{1}{(1+|z|^2)^2}d\bar{z} \wedge dz \wedge \psi\right) \\ &= *(\bar{\phi} \wedge \phi \wedge \psi) \\ &= \bar{\psi}. \end{aligned}$$

Notice that if it were true that $v \in L^2_{0,2}(\mathbb{O})$ then we would have $(\tilde{f}, \bar{\psi})_X = (f, \vartheta v)_{\mathbb{O}} = (\bar{\partial}f, v)_{\mathbb{O}} = 0$, where the boundary term in the integration by parts vanishes since f is in the domain of $\bar{\partial}_c$. Since v is not in L^2 , we need a more delicate argument to prove the claim.

Since f satisfies $\bar{\partial}_c f = 0$ in \mathbb{O} , using Friedrichs' lemma we can approximate f by a sequence $\{f_\nu\}$ of smooth forms with compact support in \mathbb{O} such that $f_\nu \rightarrow f$ in $L^2_{0,1}(\mathbb{O})$ and $\bar{\partial}f_\nu \rightarrow \bar{\partial}f = 0$ in $L^2_{0,2}(\mathbb{O})$ as $\nu \rightarrow \infty$. Thus we have

$$\begin{aligned} (\tilde{f}, \bar{\psi})_X &= (f, \vartheta v)_{\mathbb{O}} \\ &= \lim_{\nu \rightarrow \infty} (f_\nu, \vartheta v)_{\mathbb{O}} \\ &= \lim_{\nu \rightarrow \infty} (\bar{\partial}f_\nu, v)_{\mathbb{O}}, \end{aligned}$$

where the boundary term vanishes since f_ν has compact support in \mathbb{O} .

Assuming that $\bar{\partial} : L^2_{2,0}(\mathbb{O}) \rightarrow L^2_{2,1}(\mathbb{O})$ has closed range, this implies by duality that $\bar{\partial}_c : L^2_{0,1}(\mathbb{O}) \rightarrow L^2_{0,2}(\mathbb{O})$ has closed range (see Theorem 3 in [6]). Since $\bar{\partial}_c f_\nu = \bar{\partial}f_\nu$ is in the range of $\bar{\partial}_c$, there exists u_c^ν with $u_c^\nu \in \text{Dom}(\bar{\partial}_c)$ such that

$$\bar{\partial}_c u_c^\nu = \bar{\partial}f_\nu$$

and

$$\|u_c^\nu\| \leq C\|\bar{\partial}_c f_\nu\| \rightarrow 0. \quad (4.9)$$

If $\bar{\partial}$ has closed range in $L^2_{2,1}(\mathbb{O})$, then the $\bar{\partial}$ -Neumann operator $N_{2,0} : L^2_{2,0}(\mathbb{O}) \rightarrow L^2_{2,0}(\mathbb{O})$ exists. We can take $u_c^\nu = -*\bar{\partial}N_{2,0}^* \bar{\partial}f_\nu$ and u_c^ν satisfies (4.9). One can approximate u_c^ν by a sequence of smooth forms $u_c^{\nu,\epsilon}$ compactly supported in \mathbb{O} such that $u_c^{\nu,\epsilon} \rightarrow u_c^\nu$ and $\bar{\partial}u_c^{\nu,\epsilon} \rightarrow \bar{\partial}u_c^\nu$ in L^2 (see e.g. Lemma 2.4 in [19]).

Therefore

$$\begin{aligned} (\bar{\partial}f_\nu, v)_{\mathbb{O}} &= (\bar{\partial}_c u_c^\nu, v)_{\mathbb{O}} = \lim_{\epsilon \rightarrow 0} (\bar{\partial}_c u_c^{\nu,\epsilon}, v)_{\mathbb{O}} \\ &= \lim_{\epsilon \rightarrow 0} (u_c^{\nu,\epsilon}, \vartheta v)_{\mathbb{O}} \\ &= (u_c^\nu, \bar{\psi})_{\mathbb{O}} \\ &\rightarrow 0 \end{aligned}$$

where we have used (4.9). Therefore $(\tilde{f}, \bar{\psi})_X = 0$ and this shows that there is a \tilde{u} on X satisfying $\bar{\partial}\tilde{u} = \tilde{f}$. Since $\tilde{f} \in L^2(X)$, by interior elliptic regularity, we have $\tilde{u} \in W^1(X)$. Since $\tilde{u}|_{X \setminus \mathbb{O}}$ is holomorphic, it follows from Lemma 4.2, parts 2 and 3 that \tilde{u} reduces to a constant c on $X \setminus \mathbb{O}$. Therefore $u = \tilde{u} - c$ is a function

in X satisfying $\bar{\partial}u = f$ in X . The support of u is in $\bar{\mathbb{O}}$. The boundary $b\mathbb{O}$ is smooth. Thus we have, using the weak and strong extension for $\bar{\partial}_c$ (cf. [19, Lemma 2.4]) that $u \in \text{Dom}(\bar{\partial}_c)$ and $\bar{\partial}_c u = f$. \square

End of Proof of Theorem 1.1. Suppose that the range of $\bar{\partial} : L^2_{2,0}(\mathbb{O}) \rightarrow L^2_{2,1}(\mathbb{O})$ is closed in $L^2_{2,1}(\mathbb{O})$. Then by Lemma 4.4 we have $H_{c,L^2}^{0,1}(\mathbb{O}) = 0$. But this is in contradiction with the conclusion of Proposition 3.1. \square

5. FORMS ON \mathbb{O} DEPENDING ON THE \mathbb{C}^* FACTOR ONLY

Since \mathbb{O} is biholomorphic to $\mathbb{C}^* \times \mathbb{A}$, it may seem that the lack of closed range proved above might somehow be related to the factor \mathbb{C}^* . To investigate this, we consider forms on \mathbb{O} , whose coefficients depend only on the variable Z ranging over \mathbb{C}^* and not on the variable W ranging over \mathbb{A} . In the natural coordinates (z, w) of $\mathbb{P} \times \mathbb{E}$, this corresponds to considering forms on \mathbb{O} whose coefficients are functions of the inhomogeneous coordinate z on $\mathbb{P} \setminus \{\infty\}$, since $z = Z$ by (4.3).

Let us say that a $(2, 0)$ -form $u \cdot dz \wedge \psi$ on \mathbb{O} *depends on z only* if the coefficient function u is a function of the coordinate z alone. Denote by $L^2_{2,0}(\mathbb{O}, z\text{-only})$ the space of square-integrable $(2, 0)$ -forms which depend on z only. Similarly, let $L^2_{2,1}(\mathbb{O}, z\text{-only})$ be the subspace of $L^2_{2,1}(\mathbb{O})$ consisting of forms of the type $v \cdot d\bar{z} \wedge dz \wedge \psi$, with v a function of z . Note that $L^2_{2,0}(\mathbb{O}, z\text{-only})$ and $L^2_{2,1}(\mathbb{O}, z\text{-only})$ are closed subspaces of $L^2_{2,0}(\mathbb{O})$ and $L^2_{2,1}(\mathbb{O})$ respectively. Note further that each form in $L^2_{2,1}(\mathbb{O}, z\text{-only})$ is automatically closed, and the image under $\bar{\partial}$ of $L^2_{2,0}(\mathbb{O}, z\text{-only})$ is contained in $L^2_{2,1}(\mathbb{O}, z\text{-only})$. Let $T : L^2(\mathbb{C}) \dashrightarrow L^2(\mathbb{C})$ be the unbounded closed densely defined operator given on smooth functions by

$$Tu(z) = (1 + |z|^2) \frac{\partial u}{\partial \bar{z}}, \quad (5.1)$$

and with the weak maximal realization. Then we have the following:

Lemma 5.1. *There are Hilbert-space isomorphisms $\eta : L^2_{2,0}(\mathbb{O}, z\text{-only}) \rightarrow L^2(\mathbb{C})$ and $\theta : L^2_{2,1}(\mathbb{O}, z\text{-only}) \rightarrow L^2(\mathbb{C})$ such that the following diagram commutes:*

$$\begin{array}{ccc} L^2_{2,0}(\mathbb{O}, z\text{-only}) \cap \text{Dom}(\bar{\partial}) & \xrightarrow{\bar{\partial}} & L^2_{2,1}(\mathbb{O}, z\text{-only}) \\ \eta \downarrow & & \theta \downarrow \\ \text{Dom}(T) & \xrightarrow{T} & L^2(\mathbb{C}) \end{array}$$

where $\text{Dom}(T) \subset L^2(\mathbb{C})$ is the domain of the maximally realized operator T , and the map $\eta : \text{dom}(\bar{\partial}) \cap L^2_{2,0}(\mathbb{O}, z\text{-only}) \rightarrow \text{Dom}(T)$ is a bijection.

Proof. Let $\eta : L^2_{2,0}(\mathbb{O}, z\text{-only}) \rightarrow L^2(\mathbb{C})$ be the map

$$\eta : u \cdot dz \wedge \psi \mapsto u$$

and let $\theta : L^2_{2,1}(\mathbb{O}, z\text{-only}) \rightarrow L^2(\mathbb{C})$ be the map

$$\theta : v \cdot \bar{\phi} \wedge \phi \wedge \psi \mapsto \frac{v}{(1 + |z|^2)}.$$

(Note that each element of $L^2_{2,1}(\mathbb{O}, z\text{-only})$ may be written in the form $v \cdot \bar{\phi} \wedge \phi \wedge \psi$, for a function v of $z \in \mathbb{C}^*$). We first check that η and θ are isomorphisms. The space $L^2_{2,0}(\mathbb{O})$ is defined independently of the

metric, since its elements are sections of the canonical bundle, and we have

$$\begin{aligned}\|f\|^2 &= \int_{\mathbb{O}} f \wedge \bar{f} \\ &= \int_{\mathbb{O}} |u|^2 dz \wedge \psi \wedge d\bar{z} \wedge \bar{\psi} \\ &= 4 \int_{\mathbb{O}} |u|^2 (1 + |z|^2)^2 d\mathbb{O},\end{aligned}$$

where $d\mathbb{O} = (-2i)^{-2} \phi \wedge \bar{\phi} \wedge \psi \wedge \bar{\psi}$ is the Riemannian volume form of \mathbb{O} . Pulling back the integral to $\mathbb{C}^* \times \mathbb{A}$ via the map Φ of (4.3), and noting that under the pullback metric, the volume form is the form ω of (4.8), we see that

$$\begin{aligned}\|f\|^2 &= 4 \int_{\mathbb{C}^* \times \mathbb{A}} |u|^2 (1 + |Z|^2)^2 \omega \\ &= 4\beta^2 \int_{\mathbb{C}^* \times \mathbb{A}} \frac{|u|^2}{|W|^2} dV,\end{aligned}$$

where dV denotes the standard Euclidean volume form of $\mathbb{C}^2 = \mathbb{R}^4$. Therefore, if $f \in L^2_{2,0}(\mathbb{O}, z\text{-only})$, it follows by Fubini's theorem that

$$\|f\|_{L^2_{2,0}(\mathbb{O}, z\text{-only})}^2 = C \|u\|_{L^2(\mathbb{C})}^2,$$

where C is a constant independent of $f = u \cdot dz \wedge \psi$. Since $\eta(f) = u$, and η is clearly surjective, it follows that η is an isomorphism of Hilbert spaces.

We now consider the map $\theta : g \mapsto v \cdot (1 + |z|^2)^{-1}$, where $g = v \cdot \bar{\phi} \wedge \phi \wedge \psi$. We have

$$\begin{aligned}\|g\|_{L^2_{2,1}(\mathbb{O})}^2 &= \int_{\mathbb{O}} |v|^2 d\mathbb{O} \\ &= \beta^2 \int_{\mathbb{C}^* \times \mathbb{A}} |v(Z)|^2 \frac{dV}{(1 + |Z|^2)^2 |W|^2} \\ &= C \int_{\mathbb{C}^*} \frac{|v(Z)|^2}{(1 + |Z|^2)^2} dV(Z) \\ &= C \|\theta(g)\|_{L^2(\mathbb{C})}^2,\end{aligned}$$

so that θ is also an isomorphism, since it is clear that θ is surjective.

Now let $f \in \text{Dom}(\bar{\partial}) \cap L^2_{2,0}(\mathbb{O}, z\text{-only})$, and write $f = u \cdot dz \wedge \psi$ where u is a function of the variable z only. Then $T(\eta(f)) = (1 + |z|^2) \frac{\partial u}{\partial \bar{z}}$. But

$$\begin{aligned}\bar{\partial}f &= \frac{\partial u}{\partial \bar{z}} d\bar{z} \wedge dz \wedge \psi \\ &= (1 + |z|^2)^2 \frac{\partial u}{\partial \bar{z}} \bar{\phi} \wedge \phi \wedge \psi,\end{aligned}$$

so that $\theta(\bar{\partial}f) = (1 + |z|^2) \frac{\partial u}{\partial \bar{z}} = T(\eta f)$ and the diagram commutes. To see that η is bijective, note first that η is injective, since it is an isomorphism of Hilbert spaces. Further, it is clear that whenever $u \in \mathcal{C}^\infty(\mathbb{C}) \cap \text{Dom}(T)$, then u lies in the image of η . Now the surjectivity of η follows from the density of smooth forms in the domain of the maximal L^2 -realization of a differential operator. \square

We now analyze the behavior of the operator T :

Lemma 5.2. *The range of the operator $T : L^2(\mathbb{C}) \dashrightarrow L^2(\mathbb{C})$ of (5.1) is the orthogonal complement of the function $(1 + |z|^2)^{-1}$ in $L^2(\mathbb{C})$. In particular the range is closed.*

Proof. We have the decomposition $L^2(\mathbb{C}) = \overline{\text{img}(T)} \oplus \ker(T^*)$, where T^* denotes the Hilbert space adjoint of T . Integration by parts shows that for $v \in \text{Dom}(T^*)$, we have $T^*v = -\frac{\partial}{\partial z} \left\{ (1 + |z|^2)v \right\}$. Consequently, the kernel of T^* consists of $v \in L^2(\mathbb{C})$ such that $(1 + |z|^2)\bar{v}$ is entire. Denoting this entire function by h , we have $\bar{v} = (1 + |z|^2)^{-1}h$ is in $L^2(\mathbb{C})$. Expanding h in an entire Taylor series, we see that $\bar{v} \in L^2(\mathbb{C})$ if and only if h is reduced to a constant. Consequently, the space $\ker(T^*)$ is one-dimensional and spanned by $(1 + |z|^2)^{-1} \in L^2(\mathbb{C})$. Therefore

$$\overline{\text{img}(T)} = \left\{ f \in L^2(\mathbb{C}) : \int_{\mathbb{C}} f \cdot (1 + |z|^2)^{-1} dA = 0 \right\},$$

where dA is the Lebesgue measure on \mathbb{C} . Note also that T is injective from $L^2(\mathbb{C})$ to $\text{img}(T)$, since its null-space consists of L^2 holomorphic functions on \mathbb{C} , the only instance of which is the zero function.

Let now $f \in \overline{\text{img}(T)}$, i.e. the inner product $(f, (1 + |z|^2)^{-1})$ vanishes. Noting that $f \in L^2(\mathbb{C})$, we easily conclude using the Cauchy-Schwarz inequality that $f \cdot (1 + |z|^2)^{-1} \in L^1(\mathbb{C}) \cap L^2(\mathbb{C})$. Define a function γ on \mathbb{C} by $\gamma = \mathfrak{F} \left(f \cdot (1 + |z|^2)^{-1} \right)$, where \mathfrak{F} is the Fourier transform on \mathbb{R}^2 . By standard properties of the Fourier transform, the function γ is continuous, vanishes at infinity and is also in $L^2(\mathbb{C})$. Furthermore, we have $(1 + \Delta)\gamma = \hat{f} \in L^2(\mathbb{C})$, where Δ is the Laplacian, the operator with symbol $|z|^2$. Combining with the fact that $f \in L^2(\mathbb{C})$, we conclude that $\Delta\gamma \in L^2(\mathbb{C})$, so that we have $\gamma \in W_{\text{loc}}^2(\mathbb{C})$. By Sobolev embedding then, γ belongs locally to the Hölder space C^α for any $0 < \alpha < 1$. Also by definition of γ

$$\begin{aligned} \gamma(0) &= \int_{\mathbb{C}} (1 + |z|^2)^{-1} f(z) e^{-i\langle 0, z \rangle} dA(z) \\ &= (f, (1 + |z|^2)^{-1}) \\ &= 0. \end{aligned}$$

It follows therefore that for each α with $0 < \alpha < 1$, there is a C_α such that

$$|\gamma(\zeta)| \leq C_\alpha |\zeta|^\alpha.$$

Let $u = \mathfrak{F}^{-1} \left(\frac{2\gamma}{\zeta} \right)$. We assume that u is in $L^2(\mathbb{C})$ first. Then, noting that the Fourier multiplier corresponding to the operator $\frac{\partial}{\partial \bar{z}}$ is $\frac{\zeta}{2}$, we have

$$\begin{aligned} Tu &= (1 + |z|^2) \frac{\partial}{\partial \bar{z}} \mathfrak{F}^{-1} \left(\frac{2\gamma}{\zeta} \right) \\ &= (1 + |z|^2) \mathfrak{F}^{-1}(\gamma) \\ &= (1 + |z|^2) \cdot (1 + |z|^2)^{-1} f \\ &= f, \end{aligned}$$

so that u satisfies $Tu = f$. To complete the proof it suffices to show that u is in fact in $L^2(\mathbb{C})$, so that we will have $u \in \text{Dom}(T)$ and consequently $f \in \text{img}(T)$. We have for some constants $C > 0$

$$\begin{aligned} \left\| \frac{\gamma}{\zeta} \right\|_{L^2(\mathbb{C})}^2 &\leq C \int_{\zeta \in \mathbb{C}, |\zeta| \leq 1} \frac{|\zeta|^{2\alpha}}{|\zeta|^2} dA(\zeta) + \int_{\zeta \in \mathbb{C}, |\zeta| \geq 1} \frac{|\gamma|^2}{|\zeta|^2} dA(\zeta) \\ &\leq C \left(\int_{r=0}^1 r^{2\alpha-1} dr + 1 \right) \\ &< \infty, \end{aligned}$$

where A is a positive real number. This shows that $u = \mathfrak{F}^{-1} \left(\frac{2\gamma}{\zeta} \right)$ is in $L^2(\mathbb{C})$. This completes the proof. \square

Then we have the following:

Proposition 5.3. *The operator $\bar{\partial} : L^2_{2,0}(\mathbb{O}, z\text{-only}) \dashrightarrow L^2_{2,1}(\mathbb{O})$ has closed range. The range is the orthogonal complement of the form $\bar{\phi} \wedge \phi \wedge \psi$ in the closed subspace $L^2_{2,1}(\mathbb{O}, z\text{-only})$.*

Proof. According to Lemma 5.1, the range of the operator $\bar{\partial} : L^2_{2,0}(\mathbb{O}, z\text{-only}) \dashrightarrow L^2_{2,1}(\mathbb{O}, z\text{-only})$ is equal to $\theta^{-1} \circ T \circ \eta$, and since η is surjective onto $\text{Dom}(T)$, the range of $\bar{\partial}$ coincides with $\theta^{-1}(\text{range}(T))$. Since θ is an isomorphism and an isometry (up to a constant) of Hilbert spaces, and the range of T is the orthogonal complement of $v = (1 + |z|^2)^{-1}$ in $L^2(\mathbb{C})$, it follows that the range of $\bar{\partial}$ is the orthogonal complement in $L^2_{2,1}(\mathbb{O}, z\text{-only})$ of $\theta^{-1}(v) = \bar{\phi} \wedge \phi \wedge \psi$. \square

The result of Proposition 5.3 can be strengthened to prove the following:

Proposition 5.4. *The form $g = \bar{\phi} \wedge \phi \wedge \psi$ is not in the range of $\bar{\partial} : L^2_{2,0}(\mathbb{O}) \rightarrow L^2_{2,1}(\mathbb{O})$.*

Proof. In view of the previous proposition, it is sufficient to prove that if there is a solution μ in $L^2_{2,0}(\mathbb{O})$ satisfying the equation

$$\bar{\partial}\mu = g = \bar{\phi} \wedge \phi \wedge \psi$$

then there is also a solution $\lambda \in L^2_{2,0}(\mathbb{O}, z\text{-only})$.

We work in the product model $\mathbb{C}^* \times \mathbb{A}$ with natural coordinates (Z, W) . Then we have

$$g = \frac{d\bar{Z} \wedge dZ \wedge dW}{(1 + |Z|^2)^2 W},$$

and suppose that in these coordinates the solution μ is given as $\mu = u(Z, W) \frac{dZ}{1 + |Z|^2} \wedge \frac{dW}{W}$, where from $\bar{\partial}\mu = g$ we conclude that

$$\frac{\partial}{\partial \bar{Z}} \left(\frac{u(Z, W)}{1 + |Z|^2} \right) = \frac{1}{(1 + |Z|^2)^2},$$

from which we get the relation

$$u_{\bar{Z}} \cdot (1 + |Z|^2) - Z \cdot u = 1.$$

Since μ is in L^2 , it follows that

$$\int_{\mathbb{C}^* \times \mathbb{A}} \frac{|u(Z, W)|^2}{(1 + |Z|^2)^2 |W|^2} dV < \infty,$$

where dV is the volume form of \mathbb{C}^2 . Let v be defined as a function of Z by

$$v(Z) = \frac{1}{|\mathbb{A}|} \int_{\mathbb{A}} u(Z, W) dA(W)$$

where dA is Lebesgue measure on the plane. Note that

$$\begin{aligned} \frac{\partial}{\partial \bar{Z}} \left(\frac{v(Z)}{1 + |Z|^2} \right) &= \frac{v_{\bar{Z}} \cdot (1 + |Z|^2) - Z \cdot v}{(1 + |Z|^2)^2} \\ &= \frac{1}{(1 + |Z|^2)^2} \cdot \frac{1}{|\mathbb{A}|} \cdot \int_{\mathbb{A}} \left(u_{\bar{Z}} \cdot (1 + |Z|^2) - Z \cdot u \right) dA \\ &= \frac{1}{(1 + |Z|^2)^2}. \end{aligned}$$

By the Cauchy-Schwarz inequality, $|v(Z)|^2 \leq |\mathbb{A}|^{-1} \int_{\mathbb{A}} |u(Z, W)|^2 dA(W)$. Integrating this over \mathbb{C}^* it follows that $v \cdot (1 + |Z|^2)^{-1} \in L^2(\mathbb{C})$.

We set

$$\lambda = v(Z) \frac{dZ}{1 + |Z|^2} \wedge \frac{dW}{W},$$

which is a (2,0)-form. From the fact that $v \cdot (1 + |Z|^2)^{-1} \in L^2(\mathbb{C})$, it follows that $\lambda \in L^2_{2,0}(\mathbb{O}, z\text{-only})$. Further,

$$\bar{\partial}\lambda = \frac{\partial}{\partial \bar{Z}} \left(\frac{v(Z)}{1 + |Z|^2} \right) d\bar{Z} \wedge dZ \wedge \frac{dW}{W} = g.$$

The claim is proved. \square

Remarks: (1) In spite of Theorem 1.1 and Proposition 5.3, we do not have a full understanding of the L^2 theory of the manifold \mathbb{O} . In particular, it would be interesting to compute the L^2 -cohomology of \mathbb{O} in other degrees, and to know whether the form g of the preceding lemma lies in the closure of the range of $\bar{\partial}$. In the earlier examples by Serre [25] or Malgrange [22], the groups $H^{2,1}$, $H^{1,1}$ and $H^{0,1}$ are isomorphic. (2) We note that for a bounded pseudoconvex domain with smooth boundary in \mathbb{P}^n , there does not exist any hypersurface which satisfies the condition (a) in Proposition 3.1 (see [10, Theorem 1.3]).

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